## 2D Potential Problems with Periodic Boundary Conditions

We will explain here how the codes in Chapter 1 of the book "A Beginner's Course in Boundary Element Methods" can be modified to solve a particular 2D potential problem with periodic boundary conditions. The particular problem requires solving

$$
\begin{equation*}
\frac{\partial^{2} \phi}{\partial x^{2}}+\frac{\partial^{2} \phi}{\partial y^{2}}=0 \text { for }-\infty<x<\infty, 0<y<b \tag{1}
\end{equation*}
$$

subject to the boundary conditions

$$
\begin{align*}
& \phi(x, 0)=f(x) \text { for }-\infty<x<\infty \\
& \left.\frac{\partial \phi}{\partial n}\right|_{y=b}=g(x) \text { for }-\infty<x<\infty \tag{2}
\end{align*}
$$

where $f$ and $g$ are given functions which are periodic with period $a>0$, that is, they satisfy the periodic conditions

$$
\begin{equation*}
f(x+a)=f(x) \text { and } g(x+a)=g(x) \tag{3}
\end{equation*}
$$

The problem defined by (1)-(3) may be reformulated as one which requires solving

$$
\begin{equation*}
\frac{\partial^{2} \phi}{\partial x^{2}}+\frac{\partial^{2} \phi}{\partial y^{2}}=0 \text { for } 0<x<a, 0<y<b \tag{4}
\end{equation*}
$$

subject to

$$
\begin{align*}
\phi(x, 0) & =f(x) \text { for } 0<x<a \\
\left.\frac{\partial \phi}{\partial n}\right|_{y=b} & =g(x) \text { for } 0<x<a \\
\phi(0, y) & =\phi(a, y) \text { for } 0<y<b \\
\left.\frac{\partial \phi}{\partial n}\right|_{x=0} & =-\left.\frac{\partial \phi}{\partial n}\right|_{x=a} \text { for } 0<y<b \tag{5}
\end{align*}
$$

There are several different ways to modify the codes in Chapter 1 of the book to solve the problem defined by (4)-(5). One possible way is as explained below.

Discretize each of the vertical sides of the rectangular solution domain into $N_{1}$ elements and each of the horizontal sides into $N_{2}$ elements by placing
$2 N_{1}+2 N_{2}$ points in counterclockwise fashion on the boundary as explained in Chapter 1. We have a total of $2 N_{1}+2 N_{2}$ elements, that is, $N=2 N_{1}+2 N_{2}$. The first $N_{1}$ elements denoted by $C^{(1)}, C^{(2)}, \cdots, C^{\left(N_{1}-1\right)}$ and $C^{\left(N_{1}\right)}$ elements lie on the left vertical side of the rectangular domain, that is, the side, $x=0$, $0<y<b$. The elements $C^{\left(N_{1}+1\right)}, C^{\left(N_{1}+2\right)}, \cdots, C^{\left(N_{1}+N_{2}-1\right)}$ and $C^{\left(N_{1}+N_{2}\right)}$ lie on the bottom horizontal side $y=0,0<x<a ; C^{\left(N_{1}+N_{2}+1\right)}, C^{\left(N_{1}+N_{2}+2\right)}$, $\cdots, C^{\left(2 N_{1}+N_{2}-1\right)}$ and $C^{\left(2 N_{1}+N_{2}\right)}$ lie on the right vertical side $x=a, 0<$ $y<b ; C^{\left(2 N_{1}+N_{2}+1\right)}, C^{\left(2 N_{1}+N_{2}+2\right)}, \cdots, C^{\left(2 N_{1}+2 N_{2}-1\right)}$ and $C^{\left(2 N_{1}+2 N_{2}\right)}$ on the top horizontal side $y=b, 0<x<a$. With such a discretization, we see that the last two lines of (5) give

$$
\left.\begin{array}{l}
\bar{\phi}^{(k)}-\bar{\phi}^{\left(2 N_{1}+N_{2}+1-k\right)}=0  \tag{6}\\
\bar{p}^{(k)}+\bar{p}^{\left(2 N_{1}+N_{2}+1-k\right)}=0
\end{array}\right\} \text { for } k=1,2, \cdots, N_{1}
$$

If $C^{(k)}$ is a horizontal element and $\phi$ is specified on $C^{(k)}$, we give $\operatorname{BCT}(\mathrm{k})$ the value 0 and $\operatorname{BCV}(\mathrm{k})$ the specified value of $\phi$. If $C^{(k)}$ is a horizontal element and $\partial \phi / \partial n$ is specified on $C^{(k)}$, we give $\operatorname{BCT}(\mathrm{k})$ the value 1 and $\operatorname{BCV}(\mathrm{k})$ the specified value of $\partial \phi / \partial n$. If $C^{(k)}$ is a vertical elements then both $\bar{\phi}^{(k)}$ and $\bar{p}^{(k)}$ are unknowns. If $C^{(k)}$ is a vertical element, we give $\operatorname{BCT}(\mathrm{k})$ the value of 2 (the value of $\operatorname{BCV}(\mathrm{k})$ is not important here, as it is not needed in the code).

We have $4 N_{1}+2 N_{2}$ unknowns in our formulation. The boundary integral equation for the 2D Laplace's equation together with the boundary conditions on the horizontal sides and (6) can be used to generate $4 N_{1}+2 N_{2}$ linear algebraic equations. If the unknowns are $Z^{(1)}, Z^{(2)}, \cdots, Z^{\left(4 N_{1}+2 N_{2}-1\right)}$ and $Z^{\left(4 N_{1}+2 N_{2}\right)}$ then we choose $Z^{(k)}=\bar{\phi}^{(k)}$ for $k=1,2, \cdots, N_{1}, Z^{\left(N_{1}+k\right)}=\bar{p}^{(k)}$ for $k=1,2, \cdots, N_{2}, Z^{\left(N_{1}+N_{2}+k\right)}=\bar{\phi}^{(k)}$ for $k=1,2, \cdots, N_{1}, Z^{\left(2 N_{1}+N_{2}+k\right)}=$ $\bar{\phi}^{(k)}$ for $k=1,2, \cdots, N_{2}, Z^{\left(2 N_{1}+2 N_{2}+k\right)}=\bar{p}^{(k)}$ for $k=1,2, \cdots, N_{1}$, and $Z^{\left(2 N_{1}+2 N_{2}+2 N_{1}+1-k\right)}=\bar{p}^{\left(2 N_{1}+N_{2}+1-k\right)}$ for $k=1,2, \cdots, N_{1}$. The subroutine CELAP1 has to modified accordingly. It is modifed and renamed CELAPPER as listed in the file CELAPPER.for.

For a particular example to test the code, take $b=0.50, f(x)=\sin (x)$ and $g(x)=\exp (0.50) \sin (x)$. The functions $f(x)$ and $g(x)$ are periodic with period $a=2 \pi$. The main program for setting up this problem is in the file EXPRD.for. In the main program, the subroutine CELAPPER is called. Other subprograms required to run the program are CELAP2, CPF and SOLVER (together with DAXPY, DSCAL and IDAMAX) (all listed in the book). CELAP2 is called to compute the solution at any selected point in the domain $0<x<$
$2 \pi, 0<y<0.50$. The exact value given by $\phi(x, y)=\exp (y) \sin (x)$ is also printed out by the program. At selected points, numerical results obtained using $N_{1}=10$ and $N_{2}=30$ are compared with the exact solution in the table below.

| $(x, y)$ | Numerical | Exact |
| :---: | :---: | :---: |
| $(0,1,0.2)$ | 0.1202 | 0.1219 |
| $(3.0,0.4)$ | 0.2077 | 0.2105 |
| $(5.0,0.01)$ | -0.9748 | -0.9686 |
| $(6.0,0.45)$ | -0.4382 | -0.4382 |

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